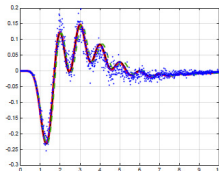
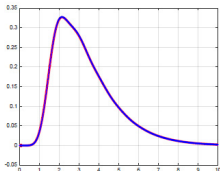


Corrections to scaled limits in random matrix theory

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- ▶ Riemann zero data and Painlevé transcendents
- ▶ Finite size corrections to Riemann zero data
- ▶ Finite size corrections in RMT scaling limits



Gap probabilities and Painléve transcendents

Painléve transcendents — solutions of one of six nonlinear 2nd order differential equations, each with the property that no moveable singular points are essential singularities.

There is a (non-autonomous) Hamiltonian theory (Malmquist). The Hamiltonian satisfies a so-called Painléve equation in sigma form, e.g. σ PII, $(\sigma'')^2 + 4\sigma'((\sigma')^2 - t\sigma' + \sigma) - a^2 = 0$.

Since the work of the Kyoto school (1980) it has been known that **gap probabilities** in random matrix theory permit evaluations in terms of sigma Painléve transcendents, e.g. (F. & Witte, 2004)

$p(s; 0; \xi)|_{\beta=2} = \frac{\pi^2}{3} s^2 \exp \int_0^{2\pi s} u(t; \xi) \frac{dt}{t}$, where

$$(su''(s))^2 + (su'(s) - u(s))(su'(s) - u(s) - 4 + 4(u'(s))^2) - 16(u'(s))^2 = 0$$

subject to the boundary condition

$$u(s; \xi) \underset{s \rightarrow 0^+}{\sim} -\frac{1}{15} s^2 + O(s^4) - \frac{\xi}{8640\pi} (s^5 + O(s^7)).$$

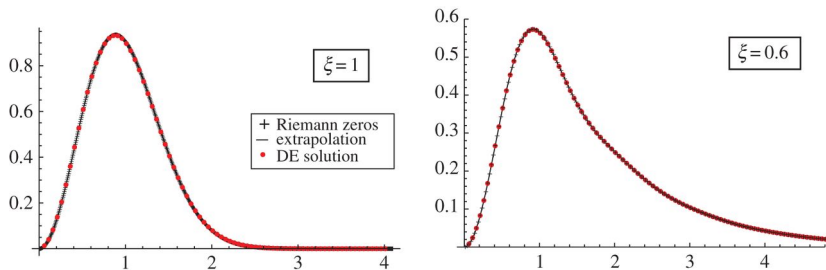
Riemann zero data

Montgomery-Odlyzko law: the statistics of the large **Riemann zeros** have the same distribution as the bulk eigenvalues of a large **complex Hermitian** random matrix.

Veracity can be probed using a high precision, big data set due to Odlyzko (~ 2000). The data set begins with zero number $10^{23} + 985,531,550$, and lists the next 10^9 .

This occurs at the point $s = 1/2 + iE$ in the complex s -plane with E equal to

$13066434408793621120027.3961465854 \approx 1.30664344 \times 10^{22}$.



A grand challenge/ suggestion — underlying determinantal point process for RZ?

Keating and Snaith (2000) hypothesised a $U(N)$ random matrix model for the leading corrections. Further developed by Bogomolny and collaborators.

The eigenvalues $e^{i\theta} = e^{i2\pi x/N}$ of $U(N)$ matrices form a determinantal point process:

$$\rho_{(k)}(x_1, \dots, x_k) = \det \left[K_N(x_i, x_j) \right]_{i,j=1}^k, \quad K_N(x, y) = \frac{1}{N} \frac{\sin \pi(x - y)}{\sin \pi(x - y)/N}.$$

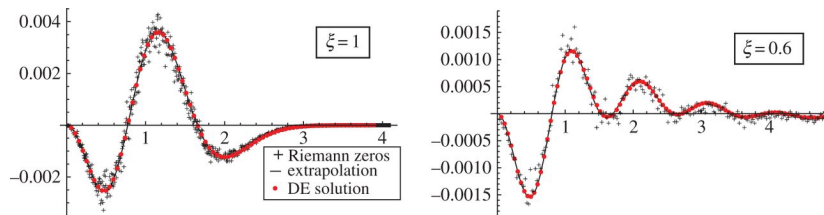
With $N = \frac{1}{\sqrt{12\Lambda}} \log \left(\frac{E}{2\pi} \right)$, $\alpha = 1 + \frac{C}{\log(E/2\pi)}$, all results known to date are consistent with

$$\rho_{(k)}^{\text{RZ}}(x_1, \dots, x_k) = \det \left[K_N^{\text{RZ}}(x_i, x_j) \right]_{i,j=1}^k$$

where

$$K_N^{\text{RZ}}(x, y) = \frac{\sin \pi(x - y)}{\pi(x - y)} + \frac{\pi(x - y) \sin(\pi(2\alpha - 1)(x - y))}{6N^2} + O\left(\frac{1}{N^4}\right).$$

The implied RMT challenge: a theory of finite size corrections



Results from F. & Witte (2004) give

$$\frac{2\pi}{N} p^N(2\pi s/N; 0; \xi) = \frac{1}{3}(N^2 - 1) \sin^2 \frac{\pi s}{N} \exp\left(-\int_0^{\pi s/N} V(\cot \phi; \xi) d\phi\right)$$

where V satisfies a particular $\tilde{\sigma}$ PVI equation. Hence

$$\begin{aligned} \frac{2\pi}{N} p^N(2\pi s/N; 0; \xi) &= \frac{\pi^2 s^2}{3} \exp\left(\int_0^{2\pi s} \frac{u^{(0)}(X)}{X} dX\right) \\ &\times \left(1 - \frac{1}{N^2} - \frac{\pi^2 s^2}{3N^2} + \frac{1}{N^2} \int_0^{2\pi s} \frac{u^{(1)}(X)}{X} dX + O\left(\frac{1}{N^4}\right)\right). \end{aligned}$$

A second order linear DE with σ PV coefficients

We find that $u^{(1)}(X)$ satisfies the second order, linear differential equation

$$\tilde{A}(s)y''(s) + \tilde{B}(s)y'(s) + \tilde{C}(s)y(s) = \tilde{D}(s),$$

where, with $u(s) = u^{(0)}(s)$, $\tilde{A}(s) = 8s^2u''(s)$. The other coefficients are also explicit polynomials in $\{u(s), u'(s), u''(s), s\}$. The equation must be solved subject to the $s \rightarrow 0^+$ boundary condition

$$u^{(1)}(s) = \frac{4}{15}s^2 - \frac{13}{6300}s^4 + \frac{\xi}{1728\pi}s^5 + O(s^6).$$

Universality results give that *all* random matrices in a certain class (for example complex Hermitian with independent entries from the same zero mean, unit variance distribution) have, after scaling, the same large N statistical properties. What about finite N corrections?

Setting up the question in the bulk

Option 1

- ▶ Choose a unitary invariant ensemble, or a complex Wigner ensemble. The simplest choice would be to consider GUE matrices.
- ▶ For complex Wigner ensembles the eigenvalue density will be to leading order given by the Wigner semi-circle law, supported to leading order on $(-\sqrt{2N}, \sqrt{2N})$. Use this density to **unfold** the eigenvalues.
- ▶ Now compute the averaged spacing distribution for some finite fraction of the eigenvalues about the origin.
- ▶ The task is to compute the large N form of this averaged spacing distribution.

Option 2

- ▶ Instead of averaging over a finite fraction of the eigenvalues, ask specifically about the spacing between, say, the two middle eigenvalues.

Gap probabilities and correlations

Let $E(n; J)$ denote the probability that the interval J contains exactly n eigenvalues. Define the generating function $G(J; \xi)$ by

$$G(J; \xi) := \sum_{n=0}^{\infty} (1 - \xi)^n E(n; J).$$

Specifically, $G(J; \xi)|_{\xi=1} = E(0; J)$.

In terms of the k -point correlations $\rho_{(k)}$,

$$G(J; \xi) = 1 + \sum_{k=1}^{\infty} \frac{(-\xi)^k}{k!} \int_{a_1}^{a_2} dx_1 \cdots \int_{a_1}^{a_2} dx_k \rho_{(k)}(x_1, \dots, x_k)$$

Remarks

- ▶ The sums terminate at N for a finite system.
- ▶ Can interpret $G(J; \xi)$ as the probability that the interval J is free of eigenvalues, in the setting that each eigenvalue has been deleted independently with probability $(1 - \xi)$.

Spacing distributions and correlations

The PDF for the event that, given there is an eigenvalue at a_1 , the next eigenvalue to the right is at a_2 , is given by

$$\begin{aligned} p((a_1, a_2)) &= -\frac{1}{\rho_{(1)}(a_1)} \frac{\partial^2}{\partial a_1 \partial a_2} E(0; (a_1, a_2)) \\ &= \frac{1}{\rho_{(1)}(a_1)} \rho_{(2)}(a_1, a_2) + \dots \end{aligned}$$

Note that this is relevant to Option 1 via the averaged quantity

$$\bar{p}(s) := \int p((a, a+s)) da = \int \frac{1}{\rho_{(1)}(a)} \rho_{(2)}(a, a+s) da + \dots$$

Difficult. Ask instead a related question: what are the properties of the finite size correction at the soft edge (i.e. neighbourhood of the largest eigenvalue).

Soft edge leading corrections

We have

$$\begin{aligned} \rho_{\max}(s) &= \frac{d}{ds} \left(\sum_{k=1}^{\infty} \frac{(-1)^k}{k!} \int_s^{\infty} dx_1 \cdots \int_s^{\infty} dx_k \rho^{(k)}(x_1, \dots, x_k) \right) \\ &= \rho_{(1)}(s) - \int_s^{\infty} \rho_{(2)}(s, x) dx + \dots \end{aligned}$$

Introduce the scaled variable

$$s_{N,t} = \begin{cases} \sqrt{2N} + t/\sqrt{2N}^{1/6}, & \text{GUE} \\ 4N + 2a + 2(2N)^{1/3}t, & \text{LUE} \end{cases}$$

F. & Trinh (2017) have shown that

$$\rho_{\max,N}(t; \xi) = \rho_{\max,\infty}(t; \xi) + \frac{1}{N^{2/3}} \tilde{\rho}(t; \xi) + O\left(\frac{1}{N}\right).$$

We know that, with $q'' = sq + 2q^3$,

$$\begin{aligned} \rho_{\max,\infty}(t; \xi) &= \frac{d}{dt} \det(\mathbb{I} - \xi \mathbb{K}_{(t,\infty)}^{\text{Airy}}) \\ &= \frac{d}{dt} \exp\left(-\int_t^{\infty} (x-t)q^2(x; \xi) dx\right) \end{aligned}$$

Weak universality?

The function $\tilde{p}(t; \xi)$ can be characterised as the solution of a 2nd order linear differential equation with σ PII coefficients. It is *different* for the GUE and LUE.

Note subtlety in relation to the Laguerre case

$s_{N,t} = 4N + 2(2N)^{1/3}t + 2a$. What would happen if instead we set $s_{N,t} = 4N + 2(2N)^{1/3}t$?

$$p_{\max,N}\left(t + \frac{c}{N^{1/3}}; \xi\right) = p_{\max,\infty}(t; \xi) + \frac{c}{N^{1/3}} \frac{d}{dt} p_{\max,\infty}(t; \xi) + O\left(\frac{1}{N^{2/3}}\right).$$

Hence, a **non-optimal** choice of scaling can be detected by the “leading” correction term being related to the scaled limit as a derivative.

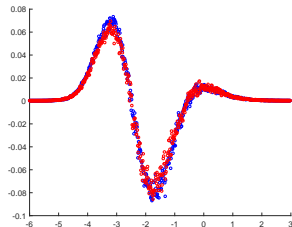
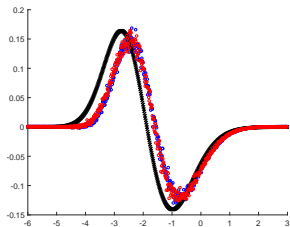
Question: Can a soft edge scaling variable always be chosen so that the optimal correction term is $O\left(\frac{1}{N^{2/3}}\right)$?

A particular Wigner ensemble

Define X so that its entries are chosen independently and uniformly from the set of four values $\frac{1}{\sqrt{2}}(\pm 1 \pm i)$.

In terms of X define the Hermitian matrix $Y = \frac{1}{2}(X + X^\dagger)$. Off diagonal entries have mean zero and variance $\frac{1}{2}$. Hence to leading order largest eigenvalue is equal to $\sqrt{2N}$.

Plot histograms associated with the scaled random variable $t = \sqrt{2}N^{1/6}(\lambda_{\max} - \sqrt{2N})$, minus $p_{\max, \infty}(t)$ multiplied by $N^{1/3}$ (left). Except for shifting by $c = 1/2$, this appears to be the graph for $\frac{d}{dt}p_{\max, \infty}(t)$. Replacing t by $t - 1/(2N^{1/3})$ leads to a $N^{-2/3}$ correction.



Future work

Study optimal soft edge scaling for the density in Gaussian and Laguerre beta ensemble with beta even. This is possible due to certain duality formulas, e.g.

$$\left\langle \prod_{l=1}^N \left(x - \sqrt{\frac{2}{\beta}} x_l \right)^n \right\rangle_{\text{GE}_{\beta, N}} = \left\langle \prod_{j=1}^n (x - ix_j)^N \right\rangle_{\text{GE}_{4/\beta, n}} .$$

Find in the Gaussian case that choosing

$$\lambda = \sqrt{2N} + \frac{1}{\sqrt{2}N^{1/6}} \left(x + \left(\frac{1}{2} - \frac{1}{\beta} \right) \frac{1}{N^{1/3}} \right)$$

gives optimal $O(N^{-2/3})$ correction. For the Laguerre case, the same (a dependent) scaling as for $\beta = 2$ is optimal.

For the unitary invariant ensemble with potential $e^{-Nx^{2m}}$ a result of Dieft and Gioev establishes m dependent values of b and γ so that $\lambda = b + \frac{x}{N^{2/3}\gamma}$ gives optimal $O(N^{-2/3})$ correction.

Presently studying similar questions at the hard edge.